

I3T: Intensity Interferometry Imaging Telescope

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ABSTRACT

We propose a new approach, based on the Hanbury Brown and Twiss intensity interferometry, to transform a Cherenkov telescope to its equivalent optical telescope. We show that, based on the use of photonics components borrowed from quantum-optical applications, we can recover spatial details of the observed source down to the diffraction limit of the Cherenkov telescope, set by its diameter at the mean wavelength of observation. For this, we propose to apply aperture synthesis techniques from pairwise and triple correlation of sub-pupil intensities, in order to reconstruct the image of a celestial source from its Fourier moduli and phase information, despite atmospheric turbulence. We examine the sensitivity of the method, i.e. limiting magnitude, and its implementation on existing or future high energy arrays of Cherenkov telescopes. We show that despite its poor optical quality compared to extremely large optical telescopes under construction, a Cherenkov telescope can provide diffraction limited imaging of celestial sources, in particular at the visible, down to violet wavelengths.

Key words: instrumentation: high angular resolution – instrumentation: interferometers – telescopes – stars: imaging.

1 INTRODUCTION

50 yr ago Antoine Labeyrie introduced the technique of speckle interferometry (Labeyrie 1970) to attain the diffraction limit of ground-based optical telescopes, overcoming the degrading effects of atmospheric turbulence that limits the image sharpness of astronomical sources to atmospheric seeing limit, i.e. 1 arcsec typically. Along with his collaborators he published a first list of stars they observed with the Palomar 5-m telescope (Gezari, Labeyrie & Stachnik 1972). Soon after Hanbury Brown and his team delivered a catalogue of 32 fundamental stellar parameters obtained with the Narrabri intensity interferometer (Hanbury Brown, Davis & Allen 1974). Intensity Interferometry is based on the Hanbury Brown and Twiss effect (HBT hereafter; Hanbury Brown & Twiss 1956) that relates the bunching of photons in the intensity correlation between two separate telescopes to the spatial Fourier transform of the star intensity distribution projected across the sky. The advent of modern amplitude interferometry (Labeyrie 1975), with its superior sensitivity over HBT obliterated efforts to construct a second generation intensity interferometer, as originally planned by Hanbury Brown (Hanbury Brown 1974). Today, optical interferometers attain hectometric baselines that reveal complex structures of stellar sources from milli-arcsecond resolutions and aperture synthesis techniques (Mourad et al. 2015; Baines et al. 2018).

On the other hand, very large Cherenkov telescope arrays, primarily built for high-energy astrophysics (Deil et al. 2008; Mazin et al. 2014) have motivated several international teams to reconsider

HBT Intensity Interferometry (HBTII) in the context of the future Cherenkov Telescope Array (CTA; Dravins et al. 2013). The future full-fledged CTA will offer several thousands of baselines up to 2.5 km and a total collecting surface of about 10 000 m². Thanks to its greatly reduced sensitivity to atmospheric degrading effects, fast single photon detectors, multispectral channels correlation, and fast digital correlators, HBTII will be a serious challenger to amplitude interferometers, specially when sub milli-arcsecond (mas) resolutions down to 10 micro-arcsec (μ as) are needed at visible wavelengths. Such resolutions can potentially open a new era of discovery, e.g. measuring the angular diameter of the brightest white dwarfs (Trippe et al. 2014), of massive early-type stars in the Magellanic Clouds, and the bright line regions (BLR) components of the closest quasars (Eracleous & Halpern 1994).

Following our successful detection of temporal and spatial bunching of stellar photons on bright sources with a 1-m telescope (Guerin et al. 2017) and with two 1-m telescopes (Guerin et al. 2018), we recently determined the geometry and distance of the luminous blue variable (LBV) super-giant P Cygni in the light of its H α emission line (Rivet et al. 2020) using a 1-nm narrow band filter. HBTII has also been successfully implemented on Imaging Atmospheric Cherenkov Telescopes (IACT) very recently on the MAGIC (Acciari et al. 2020) and VERITAS arrays (Abeyssekara et al. 2020).

The rather fast progress on the renewal of HBTII has pushed our team to consider alternative techniques and concepts for the operation of interferometric Cherenkov arrays, starting with individual Cherenkov telescopes themselves designed as intensity interferometers. The present paper demonstrates that, assuming a number of modifications, it is possible to transform a segmented single Cherenkov photon bucket into the equivalent of an optical

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telescope with an angular resolution $\sim\lambda/D$ limited by diffraction, where λ is the observing wavelength and D the telescope diameter.

In the followings we first describe the conceptual design of such an instrument: its opto-mechanical set-up, the detectors, and the digital correlator. We will develop the latter point as it has crucial importance for on- and off-line processing of data and its quality. Finally, to be meaningful in terms of astrophysical outcome, we examine the signal-to-noise ratio (SNR) aspects and the limiting magnitude. Based on this aspect we also consider image reconstruction with our proposed concept down to the diffraction limit, from modulus and closure phase quantities. We call our concept Intensity Interferometric Imaging Telescope (I3T), inspired to us by Antoine Labeyrie's comparison of stellar speckle and HBT intensity interferometries (Labeyrie 1970). I3T could be implemented on the largest operating or foreseen Cherenkov telescopes such as HESS2II and the *Large-Sized Telescopes (LST)* of the CTA array. Finally we discuss a few technical aspects of our innovative concept and conclude on how it could contribute to stellar and extra-galactic astrophysics.

2 CONVERTING A PHOTON BUCKET TO A DIFFRACTION-LIMITED OPTICAL TELESCOPE

2.1 Principle of I3T

Extremely Large Cherenkov Telescopes (ELCTs hereafter, by reference to extremely large telescopes, ELTs), such as the twin 17-m collectors of MAGIC (Mazin et al. 2014), the HESSII telescope with its $24 \times 32 \text{ m}^2$ primary mirror (Deil et al. 2008), the CTA 23-m *LSTs* (Cortina 2019), and MACE in India (Singh & Yadav 2020) made of several hundreds of elementary mirrors, also called facets, are comparable in size and collecting area to the optical ELTs under construction (Tamai et al. 2020). They benefit from fast pointing and tracking properties, but their optical quality is very poor, since their point-spread function (PSF) contains 80 per cent of encircled energy in 1 mrad. Compared to imaging optical telescopes that are limited to the seeing without adaptive optics, an ELCT is a simple photon bucket. The Cherenkov *LST* PSF is optimized to record γ -ray showers that spread over 0.2–1 degree across the IACT's field of view. It also results from pointing the optical axis of each facet towards a common focal plane that their focal spots pile up as a single PSF, comparable to a single facet spread function at best. In comparison, on-axis 1.5-m and 23-m parabolic mirrors that respect the Rayleigh optical criterion would present typically 0.33 μrad and 0.02 μrad PSFs, respectively, at 500 nm. However, both of them would be limited in practice to the atmospheric seeing of 1 arcsec, i.e. 4.8 μrad .

Transforming an ELCT to an imaging device by means of HBTII needs additional modifications and set-ups. First, the folded PSFs from the primary *facets* need to be unfolded in order to record the photon fluxes from each facet independently. This allows for pairwise and triplewise correlations computation and can be done by tip-tilting each facet, i.e. redirecting its optical axis towards a given detector across the primary focal plane array of detectors (Fig. 1). This should be straightforward, since Cherenkov telescopes are equipped by hundreds and even thousands of detectors at their primary focal plane in order to collect showers from off-axis sources with distant impact points over a large field of view, typically several degrees across the sky. However, a necessary condition is that the tip-tilt drivers mounted on the mechanical support of the primary mirror dish must have enough stroke to deviate the optical axis of a

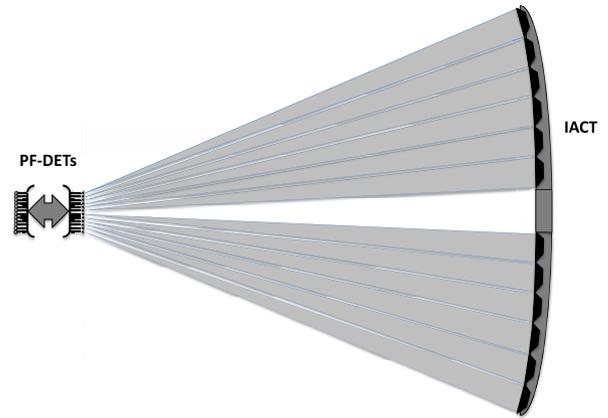


Figure 1. Layout of the optical and signal focal combiner of I3T. The spherical facet mirrors, mounted on the primary dish of a segmented Cherenkov telescope IACT, are tip-tilted so that their corresponding beams each impact a single, different detector at the primary focus of the telescope. The primary focus (PF-DETS) is equipped with as many detectors as the number of segments (facets). The individual detected photons are synchronously time-tagged and recorded in an archiving cluster of hard discs for off-line pairwise and triple-wise correlations.

given facet in the direction of a dedicated detector. At present, each such detector is 18 mm in diameter, but two neighbour detectors are 50 mm distant side to side in a hexagonal shape funnel, as specified for CTA's *LST*. Pair or triple correlation of intensities between facet mirrors provide, respectively, the Fourier power spectrum and the bi-spectrum of the source intensity across the sky. Thus an image of the observed source can be reconstructed from moduli and closure phase quantities of its Fourier transform (Malvimat, Wucknitz & Saha 2014).

2.2 Principle of intensity interferometry

In order to evaluate imaging performances of I3T we recall the principles of intensity interferometry that measures the correlation function $g^{(2)}(r, \tau)$ of the fluctuating light intensity $I(t, r)$ in the focal plane of two telescopes given by:

$$g^{(2)}(\tau, r) = \frac{\langle I(t, 0)I(t + \tau, r) \rangle}{\langle I(t, 0) \rangle \langle I(t, r) \rangle}, \quad (1)$$

where the brackets denote the average over time t and r is the baseline vector projected on the sky referenced to one of the two telescopes taken at 0 position. For a celestial source, this function is maximum at zero delay and separation. It decreases from 2 to 1 with the increasing time delay τ depending on the temporal coherence time τ_c , which is inversely proportional to the spectral bandwidth. Its dependence on the separation r between the telescopes is related to the spatial coherence of the source, i.e. the usual visibility also measured by direct stellar interferometry (Labeyrie, Lipson & Nisenson 2014), which is the normalized Fourier transform of the source brightness across the sky.

Similarly, the real term of the closure phase (Malvimat et al. 2014) can be obtained from the triple correlation $g^{(3)}(\tau, \sigma, r, s)$ (Trippie et al. 2014), defined as

$$g^{(3)}(\tau, \sigma, r, s) = \frac{\langle I(t, 0)I(t + \tau, r)I(t + \tau + \sigma, s) \rangle}{\langle I(t, 0) \rangle \langle I(t, r) \rangle \langle I(t, s) \rangle}, \quad (2)$$

where r and s are the baselines projected on the sky with two telescopes positions referenced against the third telescope and τ , σ time correlation variables.

The expressions of the SNRs for the two- and three-point intensity correlations are (Hanbury Brown 1974)

$$\text{SNR}_g^{(2)} = \alpha n_{\text{ph}} A |\gamma_{ij}|^2 \sqrt{\frac{T \Delta f}{2}}, \quad (3)$$

and (Nuñez & Domiciano de Souza 2015):

$$\text{SNR}_g^{(3)} = (\alpha n_{\text{ph}} A)^{3/2} |\gamma_{ij} \gamma_{jk} \gamma_{ki}| \Delta f \sqrt{\frac{T}{\Delta \nu}}, \quad (4)$$

where α denotes the transmission of the Cherenkov telescope including optics and detector efficiency, n_{ph} is the number of incoming photons per unit time, per unit surface, and per unit optical frequency, A is the geometrical mean of the telescope surfaces if the array is composed of elements with different diameters, T is the total observing time, Δf is the electronic bandwidth of the detection chain, and $\Delta \nu$ is the optical bandwidth. γ_{ij} , γ_{jk} , and γ_{ki} are the target's visibilities for baselines between the couples of i , j , and k telescopes and are estimated from $g^{(2)}$ and $g^{(3)}$ correlations.

These formula are valid under the following assumptions. First, the photon flux of the science target has to be much larger than the flux of spurious collected photons (e.g. from the sky background) or the dark count rate of the detector, but should not saturate the detectors. These conditions together set a working magnitude range. Another condition is that spurious correlations, for instance due to electronic noises, are negligible compared to the bunching to be measured. This condition favours the use of a narrow spectral filtering (even if $\Delta \nu$ does not appear in equation 3) in order to enhance the coherence time and thus the measured correlation (Guerin et al. 2018).

3 APPLICATION TO A 20–30 METER CLASS IACT

In order to examine the performances of the proposed I3T concept, we consider the typical case of the *LST* of CTA (Cortina 2019) that has roughly a circular shape. Implemented on a 23-m *LST*, I3T would compare, for its potential angular resolution, to a 23-m optical telescope such as the *Giant Magellan Telescope* (Martin 2019). Note that the first of the four *LSTs* is under commissioning with up to expected performances demonstrated on the Crab Nebula (CTA Observatory 2020).

3.1 *LST* primary facet configuration and baseline redundancy

The *LST* segmented primary mirror is composed of 198 hexagonal facets providing an overall collecting surface of 368 m² (Cortina 2019). Each hexagonal facet corresponds to a spherical mirror with 1.51-m long flat to flat hexagon. The spatial hexagonal paving is depicted in Fig. 2 (left), where the central facet is missing due to shadowing by the primary focus detector assembly and operational issues. Around this hole a set of six facets forms the first ring, followed outwards with 12 facets, and so on until a 23-m diameter ring. Some very external facets are absent, so that the paving does not fulfill central symmetry, rather a mirror symmetry orthogonal to the horizontal axis of rotation of *LSTs* mount.

Since the assemblage of the facets on the mechanical support of the primary mirror is hexagonal itself, there will be an important redundancy per pair and triplet combination of facets, from the closest pairs or triplets to the most distant ones. This allows adding up independent correlation measurements corresponding to the same

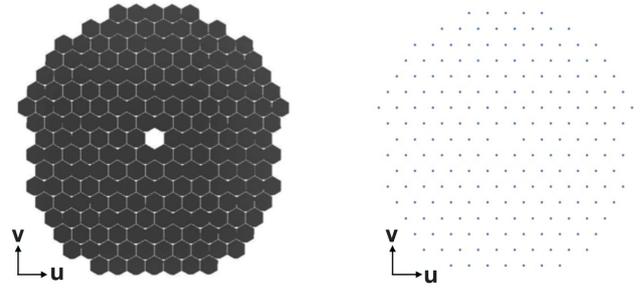


Figure 2. Left: Geometry of the facets for the *LST* (left) according to the design of CTA. Right: Its digitized representation where each facet mirror is represented by a Dirac distribution that has an amplitude proportional to the collecting area of a single facet.

baseline, increasing the corresponding SNR as the square root of the number of measurements. The number of redundancies will decrease with increasing distance between these pairs or triplets so that the SNR for measured visibilities will also decrease for increasing baselines.

There is no straightforward analytical way to determine the number of these redundancies, particularly when the facets assemblage deviates from a symmetrical distribution Hayashida 2015 around the optical axis (Fig. 2). A simple way to evaluate the redundancy is to use the modulation transfer function (MTF) of the array of 198 facets represented as the auto-correlation of a discrete set of identical Dirac distributions. Therefore, the *digitized* pupil function will result in a *digitized* MTF too, with a number of Dirac's superposed due to repeated vector baselines. The hexagonal arrangement of the 198 *LST* facet mirrors provides 382 independent visibility samples, among the 19 503 total pairwise combinations that has been directly computed from the MTF of the digitized *LST* pupil.

Let us write the digitized pupil function as follows:

$$P(u, v) = \sum_i \delta(u - u_i, v - v_i), \quad (5)$$

where u_i and v_i are the coordinates of the i -th facet with the summation carried over the 198 facets. Now, the formal MTF is the 2D auto-correlation of $P(u, v)$ that can be rewritten as an ensemble of Diracs, each one corresponding to one of the 382 independent baselines, so that each Dirac of the MTF will be multiplied by a weighting coefficient ρ_k that corresponds to the number of redundancies of the corresponding baseline:

$$\text{MTF}(u, v) = \sum_k \rho_k \delta(u - u_k, v - v_k), \quad (6)$$

where u_k and v_k are the coordinates of the spatial frequency corresponding to the baseline in the u - v plane and the summation is carried over the 382 independent spatial frequencies. Fig. 3 (top) depicts the digitized MTF. We can note that the continuous hexagonal paving of facets on the *LST* primary aperture results in a large number of identical baselines, specially for the shortest neighbouring facets in the three directions parallel to the hexagons sides. As a consequence the shortest baselines sample, the object power spectrum with the best SNR, and the redundancies will improve the measurement proportionally to $\sqrt{\rho_k}$. Apart from the zero frequency, these weights decrease from $\sqrt{180}$ at the lowest frequency of power transmissions (lowest angular resolutions), to 1 at the highest frequency (highest angular resolution). Fig. 3 (bottom) represents the azimuthal average profile of the digitized MTF and its continuous approximation ρ_c based on an arcos function depending on the radial frequency $w_c \approx$

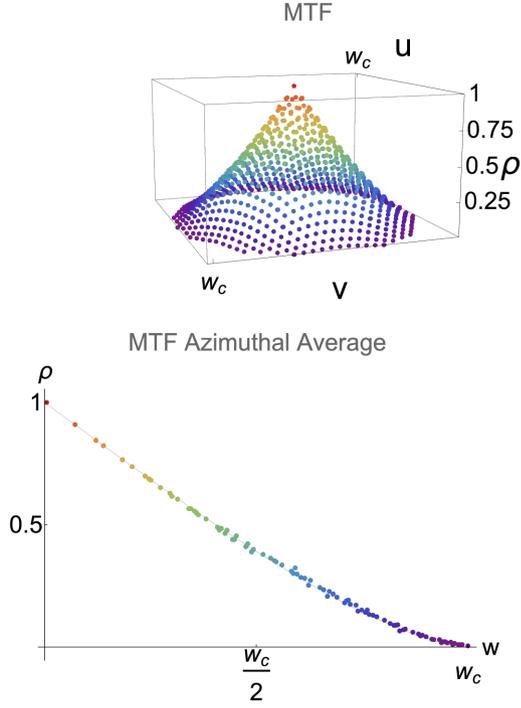


Figure 3. Top: The MTF of *LST* as the auto-correlation of its *digitized* pupil function (see Fig. 2; right), with w_c the *LST* 23-m diffraction limit cut-off frequency. Bottom: The MTF azimuthally averaged profile approximated by the following arccos-based function (see equation 7). Both graphics are normalized to the maximum of the MTF which is 198 (number of *LST* facets) with the highest redundancy of 180 at the lowest spatial frequency, i.e. the shortest baseline.

46.0 Mrad⁻¹ (Mourard et al. 1994):

$$\rho_c(w) = 2 \frac{\arccos\left(\frac{w}{23}\right) - \frac{w}{23} \sqrt{1 - \left(\frac{w}{23}\right)^2}}{\pi}, \quad (7)$$

with $w = \sqrt{u^2 + v^2}$.

3.2 SNR of the intensity correlation measurement and limiting magnitudes

Using equation (3) we can compute the SNR using the redundancy factor ρ_k of the k -th baseline, for which the visibility is γ_k , and all facets are assumed having the same collecting surface A ,

$$\text{SNR}_{g_k^{(2)}} = \alpha n_{\text{ph}} A |\gamma_k|^2 \sqrt{\frac{\rho_k T \Delta f}{2}}. \quad (8)$$

To further simplify this equation in order to estimate the limiting magnitudes of I3T, we define

$$\overline{\text{SNR}}_{g^{(2)}} = \alpha n_{\text{ph}} A \sqrt{\frac{\bar{\rho} T \Delta f}{2}}, \quad (9)$$

where γ is taken as 1 for a point-like source to estimate the limiting magnitude, $\bar{\rho} \approx 51.1$ is the value of the redundancy factor averaged over the 382 sampled spatial frequencies. In other words, it is the average power transmission of the MTF in the u - v plane. We consider a $\overline{\text{SNR}}_{g^{(2)}}$ of at least 10 to obtain quality visibility values from the observations (Trippe et al. 2014) for the image reconstruction step. Fig. 4 (top) shows the $\text{SNR}_{g^{(2)}}$ dependence with the exposure time and the apparent magnitude for a point-like source. For this

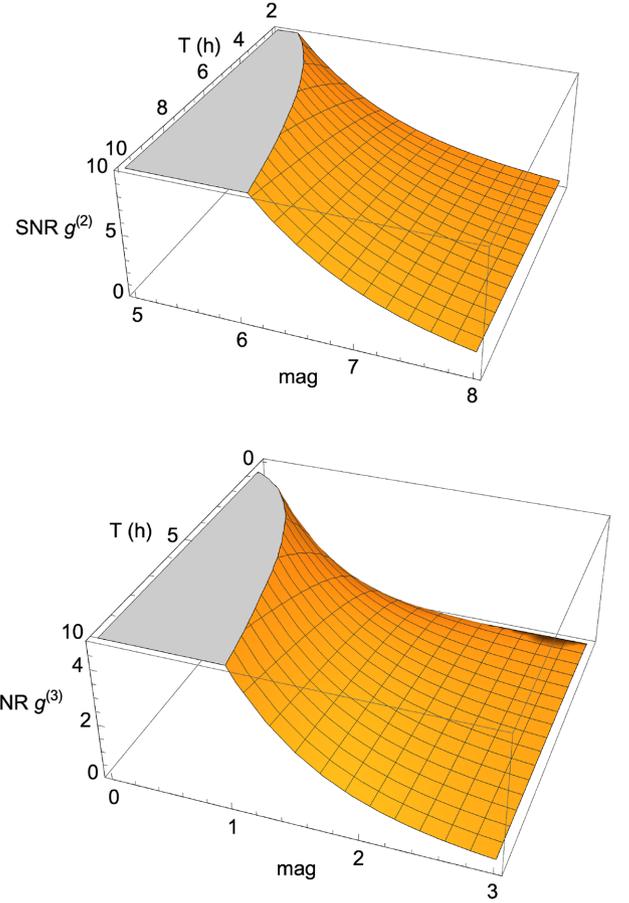


Figure 4. Top: SNR of an *LST* operated in the I3T mode for $g^{(2)}$, as a function of exposure time and magnitude. Bottom: SNR for $g^{(3)}$ as a function of the same parameters.

we also adopt the detection throughput $\alpha = 0.36$ (0.6 for the detector quantum efficiency and 0.6 for the optical throughput) and an electronic bandwidth of $\Delta f = 2$ GHz. Note that these numbers are rather optimistic compared to the presently commissioned *LST*, but could reasonably apply with the new generation of the SiPM-based detectors under development to upgrade the new *LST* camera.

We can see that for an exposure time of 3 h, an ELCT could achieve respectively a $\overline{\text{SNR}}_{g^{(2)}} \approx 5$ and $\overline{\text{SNR}}_{g^{(2)}} \approx 10$ for 6.2-mag and 5.4-mag sources. Increasing the exposure time up to 10 h the limiting magnitude improves these values to 6.85 and 6.1. Nevertheless, at the lowest reasonable $\overline{\text{SNR}}_{g^{(2)}}$ of 3, one may expect to observe still in 10 h a 7.4 magnitude point-like source. Note that for a resolved 2.5 magnitude star with a uniform disc, its second visibility maximum, i.e. 1.75 percent, could be measured with a $\overline{\text{SNR}}_{g^{(2)}}$ of 5 in 10 h.

3.3 SNR of the closure-phase measurement

So far, we discussed different aspects to determine the visibility modulus of the Fourier transform of the source brightness projected across the sky. But if we aim to reconstruct the image of that source, we also need to compute its Fourier argument components when the source is fully resolved in order to access to its spatial details. This can be achieved by computing the triplet-wise products of the visibility, namely its bi-spectrum (Hofmann & Weigelt 1987). As already mentioned (see Section 2.2), only the real term of the closure

phase can be accessed from triple correlation quantities. However, the imaginary part of the closure phase can be recovered for the image reconstruction step, thanks to the huge number of triple products and redundant baselines.

The accuracy of closure phase values will depend on the SNR for visibility triple products. In the same manner as for the pairwise SNR (see equation 3), equation 4 must be weighted by the bi-spectrum transfer function. Now, the numbers of pairwise and triplewise independent combinations for an interferometric array result from the classical $\frac{n(n-1)}{2}$ and $\frac{(n-1)(n-2)}{2}$ formula, so that the number of triple correlations, i.e. of closure phases, is 19 306 for 198 facets of *LST*, very close to the 19 503 baselines as already mentioned. We can thus compute the average number of redundant visibility phases as $\bar{\eta} \approx \frac{19306}{381} \approx 50.7$. 381 is the number of independent closure phases that we computed numerically from the triple correlation of pupil function, noting that closure phases do not provide the position of the source across the sky. Therefore, and by analogy to equation (9), the average SNR of the visibility phase to estimate the limiting magnitudes for given exposure times is

$$\overline{\text{SNR}}_g^{(3)} = (\alpha n_{\text{ph}} A)^{3/2} \Delta f \sqrt{\frac{\bar{\eta} T}{\Delta \nu}}, \quad (10)$$

Fig. 4 (bottom) depicts the dependence of the limiting magnitude of a point-like star for given $\overline{\text{SNR}}_g^{(3)}$ values as a function of the exposure time to determine the triple correlations between the facets of the *LST* telescope. Since both Δf and $\Delta \nu$ appear in $\overline{\text{SNR}}_g^{(3)}$, we take the same parameters as previously for $\overline{\text{SNR}}_g^{(2)}$ and consider a spectral bandwidth of 0.5 nm at 780 nm that sets $\Delta \nu$ to 246.4 GHz. For a first magnitude target and an integration of 10 h, the triple correlation can be determined with a $\overline{\text{SNR}}_g^{(3)}$ whilst this $\overline{\text{SNR}}_g^{(3)}$ does not exceed 1 for sources fainter than a 2.16 magnitude.

Let us consider the ratio of the SNRs for $g^{(3)}$ and $g^{(2)}$, which depends on the visibility and the quadratic ratio of Δf to $\Delta \nu$, i.e. electronic over spectral bandwidths:

$$R_{\text{SNR}32} \approx |\bar{\gamma}| \sqrt{\frac{\Delta f}{2\Delta \nu}}, \quad (11)$$

where $|\bar{\gamma}|$ is taken as the geometrical mean visibilities modulus over the three baselines of the i , j , and k facets. We see that for faint objects and/or visibilities the phase closure errors will dominate the measures and in this case phaseless image reconstruction techniques such as Fienup (1978), or more recent algorithms (Thiébaud 2008), should be preferred.

3.4 Imaging

The next step consists in setting the imaging capacity of I3T. For this, we apply a simple image reconstruction method to restore the intensity structures of a science target that would be fully resolved by a diffraction-limited optical telescope having the *LST* diameter. We take the apparent disc of a star as a relevant case, with a typical extent of 10 resolution elements (resel) across its visible hemisphere and its surface fine features. This diameter is 10 times the full width at half-maximum of an equivalent diffraction-limited telescope. A red supergiant (RSG), such as α Orionis (Betelgeuse), is a good candidate in this case since it has an angular extent of 45 mas, compared to the theoretical 4.5 mas diffraction limit of the *LST* in the visible. Also, Betelgeuse angular diameter varies significantly in R and I bands, from 600 to 800 nm wavelengths, in particular due to TiO absorption bands (Balega et al. 1982), so that the 45 mas value is taken here rather as a study case. The apparent I magnitude of Betelgeuse is -2.45 which is within the limits of SNR depicted in Fig.

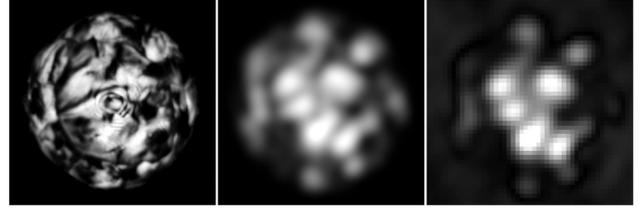


Figure 5. Image reconstruction results of a RSG star and its granulation structures following the I3T concept. Left: The RSG model (580×580 grid) with giant convective cells on its visible disc. Middle: Its theoretical image recorded by a diffraction-limited monolithic telescope with a 23 m aperture, equivalent to the *LST Cherenkov* telescope of CTA. Right: The RSG image recorded by an I3T after 10 h of noiseless observations. Note that the field rotation, i.e. azimuthal coverage in the u - v plane due to the alt-az mount of the *LST* telescope recovers partially the u - v plane gaps so that the final image corresponds fairly to the diffraction-limited image (middle).

4. Giant convective cells on the surface of Betelgeuse are inferred from spectroscopy and IR interferometry, for which intensity maps have been computed using 3D hydrodynamical simulation of stellar convection models (Chiavassa et al. 2010, 2011).

Fig. 5 depicts such a map and its image as it would be observed in the visible, say at 780 nm, by a diffraction-limited 23 m optical telescope (middle), and by a 23 m ELCT (right). The latter image is the simple inverse Fourier transform of the spatial power spectrum of the RSG, filtered by the optical transfer function of the ELCT. Here, we assume the object's Fourier power spectrum is obtained from noiseless modulus and phase data, respectively, which are determined from pairwise and triplet-wise intensity correlations obtained from facets focal plane signals. Indeed, the triple correlation of intensity signals give access only to the real part of closure phases (Rou et al. 2013; Nuñez & Domiciano de Souza 2015). However, combining phase closures equations where the huge number of redundancies compared to the small number of phase unknowns make possible to recover the imaginary parts of closure phases, so that a straightforward Fourier inversion of complex visibilities delivers the observed object's intensity across the sky.

In the real case, however, the reconstructed image will be degraded by the finite SNR of the measurements. To mimic noisy reconstructed images we apply a simplified procedure that only takes into account the $\overline{\text{SNR}}_g^{(2)}$ measurements from equation (8) that we assign as SNR for simplicity. We hamper both the real and imaginary parts of each Fourier component in the u - v plane by corresponding noise values of this SNR. Indeed this is an optimistic scenario since the determination of phase suffers from the lower $\overline{\text{SNR}}_g^{(3)}$ of equation (10).

In practice we first compute the noiseless Fourier spectrum of the RSG model that we filter by the optical transfer function of the ELCT. We thus dispose of real and imaginary parts of the visibility at u - v coordinates for $\text{MTF}(u, v) \neq 0$. Let us note γ_k the corresponding complex visibility and $\gamma_k^{(n)}$ the noisy complex value of the visibility at the k -th spatial frequency. We generate the noisy power spectrum map of the object as

$$\gamma_k^{(n)} = \gamma_k + n_k, \quad (12)$$

where n_k is a random noise that follows a normal statistical distribution according to the SNR of equation (8).

Fig. 6 depicts the reconstructed images of the RSG by Fourier inversion of visibilities for increasing values of SNRs. Here the SNR is defined as the average SNR over the measured spatial frequencies, as shown in equation (9).

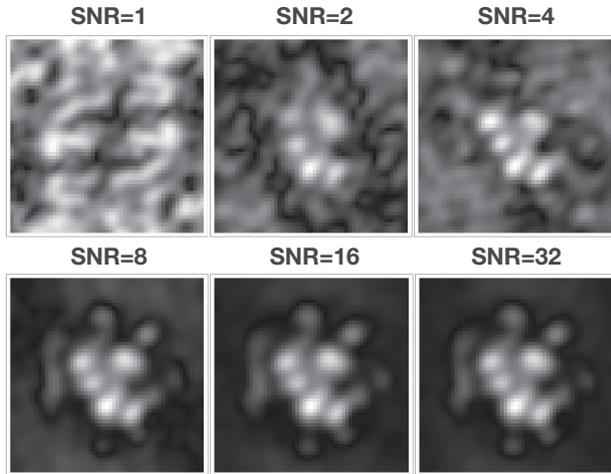


Figure 6. Left to right, top to bottom, reconstructed images of a RSG star photo-sphere for increasing SNR from 1 to 32. The adopted value of SNR is taken as its average over the measured spatial frequencies. It can be seen that for SNR = 1 the reconstructed image of the star has a random speckled structure, typically the size of the PSF, with a notable axisymmetric distribution. As the SNR increases, the reconstructed image of the convection giant cells at the surface of the RSG is getting closer to the RSG noiseless image recorded by a theoretically diffraction-limited optical telescope with the same diameter as the LST (see Fig. 5 middle).

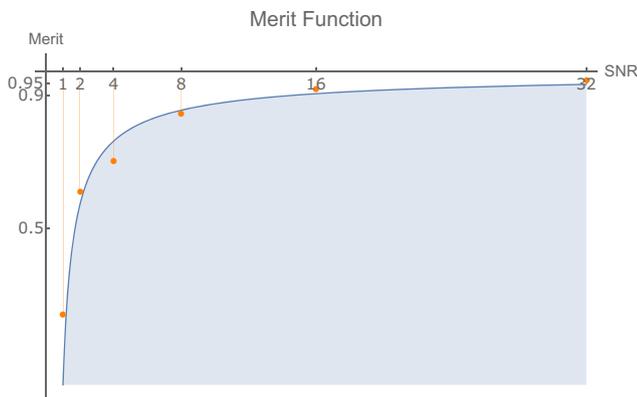


Figure 7. Evaluation of the I3T imaging capacity, based on the direct Fourier inversion reconstruction method: The MF (equation 13) is plotted for increasing values (red dots) of the average SNR of visibility modulus values (equation 9) and computed numerically as a function of SNR (see text).

The quality of the reconstructed image is already quite good for a $\text{SNR} \geq 10$, with in particular the clear identification of the giant convective cells at the surface of the RSG.

To be more quantitative, we compute the squared difference between the reconstructed image $\text{Image}(\alpha, \delta)$ and the model $\text{Model}(\alpha, \delta)$ as it would be recorded by the noiseless digitized pupil of the LST according to the I3T concept. The corresponding merit function (MF), computed as,

$$\text{MF} = 1 - \frac{\sum_{\alpha, \delta} [\text{Model}(\alpha, \delta) - \text{Image}(\alpha, \delta)]^2}{\sum_{\alpha, \delta} \text{Model}(\alpha, \delta)^2}, \quad (13)$$

is plotted as a function of the average SNR in Fig. 7. The continuous curve in this figure corresponds to $\text{MF}_{\text{SNR}} = 1 - \frac{3}{4} \text{SNR}^{-\frac{3}{4}}$ that

represents the trend of the reconstructed image quality approximated for SNR values of 1, 2, 4, 8, 16, and 32. One sees that for SNRs larger than 10, errors on the reconstructed image is about 5 per cent and better.

4 DISCUSSION

The concept of I3T (see Fig. 1) implies the tip-tilt of spherical facet mirrors on the parabolic primary mount, to have enough stroke enabling them to separately point and focus collected photons upon their dedicated detector at the common unfolded focal plane array of detectors. Now the LST has a quite open aperture, i.e. $F/D = 1.2$ for $F = 28$ m, which is also the focal length of individual facets, so that the linear size of their PSF is 28 mm for 80 per cent energy concentration (Cortina & Teshima 2016). Each facet is supported on the tetrahedral primary structure, on a triangular support, and can be oriented in tip-tilt by fine positioning actuators with a stroke of 36 mm and a $5 \mu\text{m}$ resolution. This stroke, given the focal distance of a facet, makes possible to deviate its PSF for about 50 cm from the on-axis focal point of LST therefore sufficient to record separately intensity signals for each facet according to Fig. 1 and with much better sampling of the PSF with the upgraded 7420 pixels focal camera of the LST (Alispach & et al. 2020).

Besides changing the tip-tilt setpoint of the facets, the practical implementation of I3T on an LST requires an additional component to be installed in front of the focal detectors array: a specially designed mask incorporating an array of narrow band optical filters (typically 1 to 0.1 nm), and cone-shaped baffles (one per facet). The role of the optical filters is to increase the light coherence time (see Section 2.2). The purpose of the baffles is to reduce sky background cross-contamination. Indeed, without baffles, a given focal detector receives the target's light reflected by the facet to which it corresponds, but receives also the sky background light reflected by all facets of the dish. This spurious light is non-coherent and its residual correlation can be calibrated out, but its fluctuation would hamper severely the SNR. The detailed specification of such an intrafocal component is clearly out of the scope of this article, but its presence is mandatory to achieve the limiting magnitudes set by fundamental noise and inherent detector electronics.

To improve the I3T concept sensitivity for faint science targets, the flexibility of LST primary mirror makes possible to group a number of neighbouring facets as a single superfacet, by overlapping their focal images on the top of each other on a same detector. For instance a combination of a hexagonal set of seven facets results in paving the 198 facets in a much smaller number of super-facets, but improves the SNR for any pair of superfacets by a factor of 7. The price to pay is a poorer sampling of the u-v plane, with much fewer spatial frequencies. Other combinations are also possible. One may consider this as a hierarchically digitized pupil of an ELCT, where facets grouping adapt to the complexity and magnitude of the observed source. Note also that the correlations can be added-up over as many nights as required to reach the necessary SNR and then carry the image reconstruction from the complex visibility estimates.

One also needs to consider the computational power needed for pair and triplet-wise intensity correlations between the 198 facets of an LST, which might sound considerably demanding. This turns out to be a moderate limitation of the I3T concept in fact. Present IACT intensity interferometers concentrate the photon flux from their entire aperture upon a single focal detector. Whereas in our case one can use time-tagged photon detector for recording the fluxes collected by different facets. Thus, we only need to store the arrival time of all photons, which is much less information: typically a few hundreds

of Gbytes per facet and for 10 h of integration time on a bright star, than what would produce a continuous signal digitized at a high rate.

Computing all the correlations between all pairs and triplets is certainly demanding but can be done off-line and thus doesn't represent a bottleneck.

Another advantage of computing the correlations off-line is that it allows to compensate for the field-rotation due to the alt-azimuthal mount of the LST. Indeed, the vector baselines between facet pairs change in orientation as projected on the sky, so that for faint sources demanding long exposures, one needs to correct for field rotation. The derotation can be done numerically, when computing the correlations, by adapting the facet pairs, corresponding to a given spatial frequency, as a function of the hour angle, before averaging the data over time.

5 CONCLUSION

In this paper we described a concept, called I3T, to operate a large Cherenkov telescope as an equivalent diffraction-limited optical telescope based on the technique of intensity interferometry introduced by Hanbury Brown and Twiss more than 60 yr ago. The concept can be applied to short visible wavelengths, in the violet and beyond, within atmospheric transmission limits. Foreseen extremely large telescopes can hardly operate at their diffraction limit, unless laser-guide adaptive optics remove atmospheric turbulence degrading effects, or passive techniques such as classical speckle interferometry (Labeyrie 1970) are used.

We discussed a few technical and operational issues of the I3T concept for its implementation on the first LST of the CTA telescopes, recently operated on the Canarias observatory next to the MAGIC Cherenkov array. We have shown that there is no show-stopper in principle to transform an LST unit to an imaging diffraction-limited telescope. A practical implementation of our I3T concept needs, however, to be assessed in much more technical details, but that is beyond the scope of this paper where we focused on the principles.

We remind that the full-fledged CTA array will operate with four LST IACTs. Therefore if these are configured according to the proposed I3T, the correlation signals can be processed separately and co-added to gain by a factor of four on the limiting magnitude of a single LST. Or equivalently this would divide by a factor of 4 the long exposure to attain the same SNR for a given magnitude compared to a single LST observation.

On the other hand, if the whole CTA is operated as an intensity interferometric synthetic array, and if all the individual telescopes function in the I3T mode for their facet sub-pupils, then the u-v coverage will have much better spatial frequency resolution in the Fourier plane around a given spatial frequency defined by the baseline between two LSTs. Indeed facets of LST, Medium-Sized Telescopes (MST), and Small-Sized Telescopes (SST) could be combined in pair and triplet combinations with SNRs now proportional to the geometrical mean surfaces for hybrid combinations. Hence, the u-v plane will be filled by much more resolution elements, resulting in better tracking of visibility gradients in the Fourier plane, specially around zero modulus visibilities, and more robust reconstructed image of the science object under study. Also, the I3T technique can readily be applied to existing Cherenkov telescopes like HESS28 in Namibia for instance that presently represents the largest available collector of this type: $24 \times 30 \text{ m}^2$ primary mosaic telescope composed of 850 flat-to-flat 90 cm facets therefore with better performances than the CTA LST collectors.

As for astrophysical programs, if I3T configuration is implemented on CTA, a number of targets will become accessible for a 23

m imaging ELCT in the visible wavelengths. Stellar astrophysics with surface structures imaging of AGB and post-AGB stars will become practicable with reasonable total exposure times, noting that correlation quantities can be added night to night to attain the desired SNR value necessary to restore the fine details of their structures. Massive early type stars, like LBVs, Be, and B[e] stars circumstellar envelopes, are favorite targets for an ELCT, as well as massive binary stars like Wolf-Rayet (WR) exhibiting colliding winds and emission of H, He, and CNO lines signatures. Also, pre-main sequence stars like Ae/Be Herbig or β Pictoris debris discs could be imaged with resolutions on the order of 4 mas, unravelled otherwise by existing or planned ELT's when we consider short wavelengths down to 400 nm.

Of special interest is the added value of I3T for extra-galactic studies, specially on globular clusters or starburst compact clusters like R136 in the Large Magellanic Cloud. The best resolved images of this object were obtained with *HST* and the SPHERE (Khorrami et al. 2017) on the VLT high dynamic imaging in J, H, and K bands, whereas an ELCT would offer visible imaging of this object with a spatial resolution 10 times better than the *HST*. Last but not least, the central BLR components of the closest quasars, like 3C 273, could be scrutinized in Balmer emission lines such as H δ , with roughly an exposure time of 10 h according to Fig. 4. Indeed such ambitious programs will pose the serious problem of the night sky background, which is 6–8 mag in any 1 mrad/4 arcmin diameter aperture due to IACTs very poor PSF optical quality. Solutions exist, however, to overcome this. For example, by relaying the focal image to the Cassegrain focus of the IACT, where active optics can correct for the optical aberrations up to the required level where sky background is no more a limit for extragalactic targets. The details of such solutions are beyond the scope of the present paper that need end-to-end realistic studies.

Implementing the I3T concept on large Cherenkov telescopes, such as the South and North LSTs, HESS2, MAGIC, or MACE provides a 'cheaper' solution compared to optical ELT's, to obtain diffraction-limited imaging telescopes, especially at very short visible wavelengths, with reduced sensitivity to atmospheric turbulence limitations (seeing, coherence time, and spectral dispersion).

Finally, it is noteworthy that, in their historical developments, homodyne interferometry came first, then its developments stopped, HBTII appeared and pulled ahead, then HBTII stopped and homodyne interferometry moved ahead again, and now HBTII may pull ahead again, at least for kilometric synthesis arrays in the visible wavelengths.

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DATA AVAILABILITY

No new data were generated or analysed in support of this research.

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